## Phys 487 Discussion 7 - Degenerate Perturbation Theory

Today we will learn and use the procedure for applying perturbation theory to degenerate states; in the next lecture, we will derive the procedure.

## The Procedure

Perturbation theory always starts with an "unperturbed" Hamiltonian $H_{0}$ whose eigenstates $\left\{\left|n^{(0)}\right\rangle\right.$ or $\left.\psi_{n}^{(0)}\right\}$ and eigenvalues $\left\{E_{n}^{(0)}\right\}$ can be obtained exactly. A small perturbing Hamiltonian $H^{\prime} \ll H_{0}$ is then added to $H_{0}$ to produce the full Hamiltonian $H=H_{0}+\varepsilon H^{\prime}$. This is the Hamiltonian whose eigen-things we would like to obtain. I have attached a dimensionless scale factor $\varepsilon \ll 1$ to $H^{\prime}$ so that I can easily keep track of orders of smallness. (Sometimes such a small scale factor is an intrinsic part of the problem, sometimes not.)
Suppose that a subset of the unperturbed eigenenergies $\left\{E_{n}^{(0)}\right\}$ are degenerate, i.e. have the same value $E_{\alpha}$. Let the quantum numbers of these degenerate eigenstates be $\{\alpha 1, \alpha 2, \alpha 3, \ldots, \alpha n\}$. If we write $H_{0}$ in matrix form using as basis the unperturbed eigenstates $\left\{\left|n^{(0)}\right\rangle\right\}$, we get the diagonal matrix $\left(\mathbf{H}_{0}\right)_{m n} \equiv\left\langle m^{(0)}\right| \hat{H}_{0}\left|n^{(0)}\right\rangle$ :

$$
\mathbf{H}_{0}=\left(\begin{array}{ccccc}
E_{1}^{(0)} & & & & \\
& E_{2}^{(0)} & & & \\
\\
& & \boldsymbol{E}_{\alpha} & & \\
\\
& & & \boldsymbol{E}_{\alpha} & \\
\\
& & & E_{5}^{(0)} & \\
& & & & \\
& & & & \\
& & \ldots
\end{array}\right) \text { where all the empty elements are } 0
$$

I have bold-faced the degenerate energies and left off the superscript (0) so that you can spot them easily. The degenerate states $\left\{\left|\alpha_{1}^{(0)}\right\rangle, \ldots,\left|\alpha_{n}^{(0)}\right\rangle\right\}$, which are just $\left\{\left|3^{(0)}\right\rangle,\left|4^{(0)}\right\rangle\right\}$ here, form a degenerate subspace where any linear combination of the $\mid \alpha_{i}$ >'s is also an eigenstate of $H_{0}$ with the same eigenvalue $E_{\alpha}$.

Degenerate perturbation theory is accomplished by finding a particular set of linear combinations of the $\mid \alpha_{i}>$ 's, i.e. within the degenerate subspace, that diagonalizes the perturbation matrix $\left(\mathbf{H}^{\prime}\right)_{i j} \equiv\left\langle i^{(0)}\right| \hat{H}^{\prime}\left|j^{(0)}\right\rangle$.
Once you have found these linear combinations $\left\{\left|\beta_{1}^{(0)}\right\rangle, \ldots,\left|\beta_{n}^{(0)}\right\rangle\right\}$, i.e. the eigenvectors of $H^{\prime}$ within the degenerate subspace, find their corresponding eigenvalues and you will have your first-order corrections

$$
E_{\beta_{i}}^{(1)}=\left\langle\beta_{i}^{(0)}\right| H^{\prime}\left|\beta_{i}^{(0)}\right\rangle
$$

If you all goes well, these $1^{\text {stt-order energy corrections should be non-degenerate, i.e. they should break the }}$ degeneracy of the $\alpha 1, \ldots$ an subspace.

In last week's discussion you had this Hamitonian :

$$
\mathbf{H}=V_{0}\left(\begin{array}{ccc}
(1-\varepsilon) & 0 & 0 \\
0 & 1 & \varepsilon \\
0 & \varepsilon & 2
\end{array}\right)
$$

where $V_{0}$ is a constant and $\varepsilon$ is a small number $\ll 1$. Have a look at the posted questions and solution to remind yourself: you used various techniques to find the $1^{\text {st }}$ and $2^{\text {nd }}$ order corrections to this 3 -state Hamiltonian's energies. Today, use degenerate perturbation theory to find the first-order correction to the two initially degenerate eigenvalues. Compare with your various results from last week!

## Problem 2 : Stark Effect

Griffiths 6.36
When an atom is placed in a uniform external electric field $\mathrm{E}_{\text {ext }}$, the energy levels are shifted - a phenomenon known as the Stark effect. In this problem we analyze the Stark effect for the $n=1$ and $n=2$ states of hydrogen. Let the field point in the $z$ direction, so the potential energy of the electron is

$$
H_{S}^{\prime}=e E_{\mathrm{ext}} z=e E_{\mathrm{ext}} r \cos \theta
$$

Treat this as a perturbation on the simple "Bohr" Hamiltonian for the hydrogen atom,

$$
H_{0}=-\frac{\hbar^{2} \nabla^{2}}{2 m}-\frac{e^{2}}{4 \pi \varepsilon_{0} r}
$$

Spin is irrelevant to this problem so ignore it.
(a) Show that the ground state energy is not affected by this perturbation, to first order.
(b) The first excited state is 4-fold degenerate: $\psi_{200}, \psi_{211}, \psi_{210}, \psi_{21-1}$. Using degenerate perturbation theory, determine the first-order corrections to the energy. Into how many levels does $\mathrm{E}_{2}$ split?
(c) What are the " $\beta$ " wave functions for part (b), i.e. the ones that diagonalize the perturbation $H_{S}^{\prime}$ ? (Griffiths calls these "good" wavefunctions.) Find the expectation value of the electric dipole moment ( $\vec{p}_{e}=-e \vec{r}$ ) in each of these "good" states. Notice that the results are independent of the applied field - evidently hydrogen in its first excited state can carry a permanent electric dipole moment.

HINT: There are a lot of integrals in this problem, but almost all of them are ZERO. So study each one carefully, before you do any calculations! For example, if the $\phi$ integral vanishes, there's not much point in doing the $r$ and $\theta$ integrals! Partial answer: $H_{13}^{\prime}=H_{31}^{\prime}=-3 e a_{0} E_{\text {ext }}$ where $a_{0}$ is the Bohr radius as usual; all other elements of $\mathbf{H}_{i j}^{\prime} \equiv\left\langle\beta_{i}^{(0)}\right| H^{\prime}\left|\beta_{j}^{(0)}\right\rangle$ are zero $\odot$.

