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Reading: FPE, Sections 5.1–5.4, 6.1
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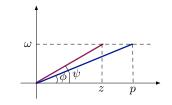
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- if z < p, then  $\psi \phi > 0$  (phase lead)
- if z > p, then  $\psi \phi < 0$  (phase lag)



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We use lag controllers as dynamic compensators for approximate PI control.

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Tracking a constant reference: assuming closed-loop stability, the FVT gives

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$$(s+p)(s-1) + K(s+z) = 0$$
$$s^2 + (K+p-1)s + Kz - p = 0$$
  
Conditions for stability:  $K \ge 1$  ,  $n \in Kz \ge n$ 

Conditions for stability: K > 1 - p, Kz > p

Tracking a constant reference: if the stability conditions

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are satisfied, then the steady-state error is

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Lag compensation *does not* give perfect tracking (indeed, it does not change system type), but we can get as good a tracking as we want by playing with K, z, p. On the other hand, unlike PI, lag compensation gives a stable controller.

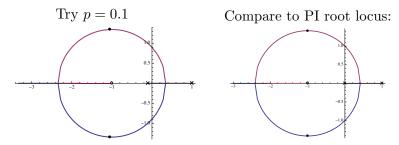
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Intuition: By choosing p very close to zero, we can make the root locus arbitrarily close to PI root locus (stable for large enough K). Let's check:

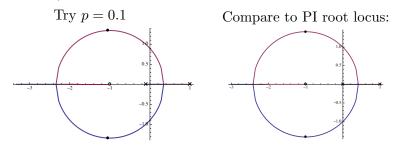
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What do we see? Compared to PD vs. lead, there is no qualitative change in the shape of RL, since we are not changing #(poles) or #(zeros).

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Main technique: select parameters to satisfy the phase condition (points on RL must be such that  $\angle L(s) = 180^{\circ}$ ).

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Main technique: select parameters to satisfy the phase condition (points on RL must be such that  $\angle L(s) = 180^{\circ}$ ).

Caveat: may not always be possible!

Pole Placement via RL

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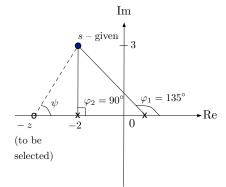
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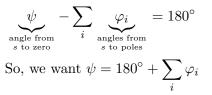
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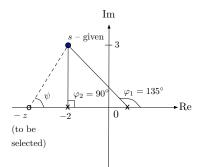
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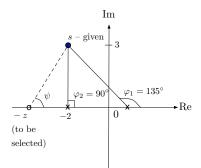
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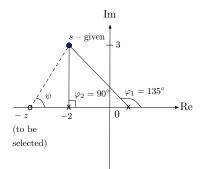
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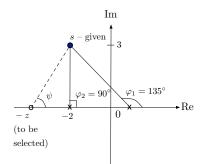




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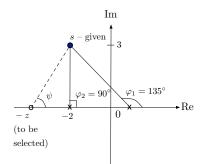
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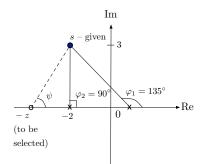


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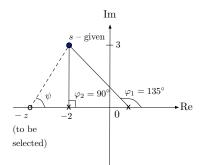


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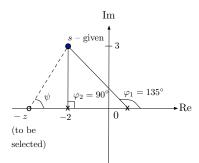
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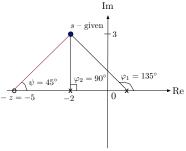
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compute s.s. tracking error:

$$\left|\frac{1}{1-\frac{Kz}{p}}\right| = \frac{1}{6.5} \approx 15\%$$

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- provides stability and perfect steady-state tracking of constant references
- ► replace unstable I-controller K/s with a stable lag controller  $K\frac{s+z}{s+p}$ , where p < z
- ▶ this does not change the shape of RL compared to PI

Obvious solution — combine lead and lag compensation.

We will develop this further in homework and later in the course using frequency-response design methods — which are the subject of several lectures, starting with today's.

### The Frequency-Response Design Method Recall the frequency-response formula:

$$\sin(\omega t) \longrightarrow G(s) \longrightarrow M \sin(\omega t + \phi)$$

where  $M = M(\omega) = |G(j\omega)|$  and  $\phi = \phi(\omega) = \angle G(j\omega)$ 

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$$\sin(\omega t) \xrightarrow{} G(s) \xrightarrow{} M \sin(\omega t + \phi)$$
  
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- 1.  $u(t) = e^{st} \longrightarrow y(t) = G(s)e^{st}$ 2. Euler's formula:  $\sin(\omega t) = \frac{e^{j\omega t} - e^{-j\omega t}}{2j}$
- 3. By linearity,

$$\sin(\omega t) \longmapsto \frac{G(j\omega)e^{j\omega t} - G(-j\omega)e^{-j\omega t}}{2j}$$

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Derivation:

1. 
$$u(t) = e^{st} \longmapsto y(t) = G(s)e^{st}$$
  
2. Euler's formula:  $\sin(\omega t) = \frac{e^{j\omega t} - e^{-j\omega t}}{2j}$ 

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$$\sin(\omega t) \longmapsto \frac{G(j\omega)e^{j\omega t} - G(-j\omega)e^{-j\omega t}}{2j} \ G(j\omega) = M(\omega)e^{j\phi(\omega)}$$

$$\sin(\omega t) \longrightarrow G(s) \longrightarrow M \sin(\omega t + \phi)$$

where  $M = M(\omega) = |G(j\omega)|$  and  $\phi = \phi(\omega) = \angle G(j\omega)$ 

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The Frequency-Response Design Method

$$\sin(\omega t) \longrightarrow G(s) \longrightarrow M \sin(\omega t + \phi)$$

where  $M = M(\omega) = |G(j\omega)|$  and  $\phi = \phi(\omega) = \angle G(j\omega)$ 

$$G(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$$

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$$= \left| \frac{1}{1 - \left(\frac{\omega}{\omega_n}\right)^2 + 2\zeta\frac{\omega}{\omega_n}j} \right|$$
$$= \frac{1}{\sqrt{\left[ 1 - \left(\frac{\omega}{\omega_n}\right)^2 \right]^2 + 4\zeta^2 \left(\frac{\omega}{\omega_n}\right)^2}}$$

# The Frequency-Response Design Method

For our prototype 2nd-order system:

0.4

0.2

0.5

1.0

$$G(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$$

$$M(\omega) = \frac{1}{\sqrt{\left[1 - \left(\frac{\omega}{\omega_n}\right)^2\right]^2 + 4\zeta^2 \left(\frac{\omega}{\omega_n}\right)^2}} = \frac{1}{\sqrt{1 + (4\zeta^2 - 2)\left(\frac{\omega}{\omega_n}\right)^2 + \left(\frac{\omega}{\omega_n}\right)^4}}$$

$$M(\omega)$$

$$1.0$$

$$0.8$$

$$0.6$$

$$-\zeta = 1/2$$

$$-\zeta = 1/\sqrt{2}$$

 $\zeta = 1$ 

2.5

2.0

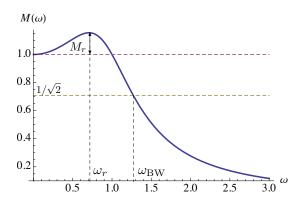
1.5

ω

 $3.0 \omega_n$ 

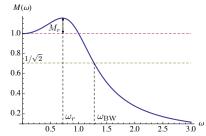
## Frequency Response Parameters

Here is a typical frequency response magnitude plot:



 $\omega_r$  – resonant frequency  $M_r$  – resonant peak  $\omega_{\rm BW}$  – bandwidth

## Frequency Response Parameters



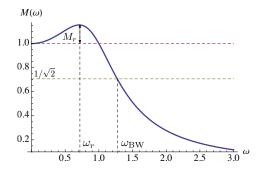
We can get the following formulas using calculus:

$$\begin{cases} \omega_r = \omega_n \sqrt{1 - 2\zeta^2} \\ M_r = \frac{1}{2\zeta\sqrt{1 - \zeta^2}} - 1 \qquad \text{(valid for } \zeta < \frac{1}{\sqrt{2}}; \text{ for } \zeta \ge \frac{1}{\sqrt{2}}, \, \omega_r = 0) \\ \omega_{\text{BW}} = \omega_n \underbrace{\sqrt{(1 - 2\zeta^2) + \sqrt{(1 - 2\zeta^2)^2 + 1}}}_{=1 \text{ for } \zeta = 1/\sqrt{2}} \end{cases}$$

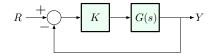
— so, if we know  $\omega_r, M_r, \omega_{BW}$ , we can determine  $\omega_n, \zeta$  and hence the time-domain specs  $(t_r, M_p, t_s)$ 

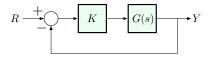
# Frequency Response & Time-Domain Specs

All information about time response is also encoded in frequency response!!

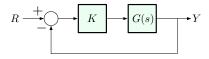


small  $M_r \longleftrightarrow$  better damping large  $\omega_{\rm BW} \longleftrightarrow$  large  $\omega_n \longleftrightarrow$  smaller  $t_r$ 



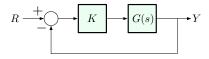


Two-step procedure:



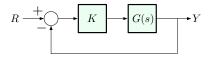
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1. Plot the frequency response of the open-loop transfer function KG(s) [or, more generally, D(s)G(s)], at  $s = j\omega$ 



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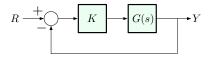
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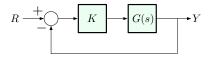


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1. Bode plots: magnitude  $|KG(j\omega)|$  and phase  $\angle KG(j\omega)$  vs. frequency  $\omega$  (could have seen it earlier, in ECE 342)



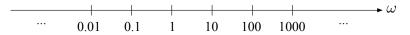
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  - 1. Bode plots: magnitude  $|KG(j\omega)|$  and phase  $\angle KG(j\omega)$  vs. frequency  $\omega$  (could have seen it earlier, in ECE 342)
  - 2. Nyquist plots:  $\operatorname{Im}(KG(j\omega))$  vs.  $\operatorname{Re}(K(j\omega))$  [Cartesian plot in *s*-plane] as  $\omega$  ranges from  $-\infty$  to  $+\infty$

#### Horizontal ( $\omega$ ) axis:

we will use *logarithmic scale* (base 10) in order to display a wide range of frequencies.

Note: we will still mark the values of  $\omega$ , not  $\log_{10} \omega$ , on the axis, but the *scale* will be logarithmic:



Equal intervals on log scale correspond to decades in frequency.

Vertical axis on magnitude plots:

we will also use logarithmic scale, just like the frequency axis.

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Reason:

$$|(M_1 e^{j\phi_1})(M_2 e^{j\phi_2})| = M_1 \cdot M_2$$
$$\log(M_1 M_2) = \log M_1 + \log M_2$$

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— this means that we can simply *add* the graphs of  $\log M_1(\omega)$ and  $\log M_2(\omega)$  to obtain the graph of  $\log (M_1(\omega)M_2(\omega))$ , and graphical addition is easy.

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Decibel scale:

$$(M)_{\rm dB} = 20 \log_{10} M$$
 (one decade =  $20 \,\mathrm{dB}$ )

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we will plot the phase on the usual (linear) scale.

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$$= \phi_1 + \phi_2$$

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— this means that we can simply add the phase plots for two transfer functions to obtain the phase plot for their product.

# Scale Convention for Bode Plots

	magnitude	phase
horizontal scale	log	log
vertical scale	log	linear

Advantage of the scale convention: we will learn to do Bode plots by starting from simple factors and then building up to general transfer functions by considering products of these simple factors.